On the relation between charge and topology

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1978 J. Phys. A: Math. Gen. 11795
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## Corrigenda

## Quantised fields over de Sitter space

Grensing G 1977 J. Phys. A: Math. Gen. 10 1687-719
On p 1688 formula (2.4) should read

$$
\begin{equation*}
g^{\dagger} E g=E \quad g^{\mathrm{T}} E^{\prime} g=E^{\prime} \tag{2.4}
\end{equation*}
$$

and on $p 1708$ formula (9.11) should read

$$
\begin{align*}
& \Delta^{-}\left(z_{1}, z_{2}\right)=\frac{1}{2} \mathrm{i}(2 \pi)^{-\frac{1}{2} d} \Gamma\left(\frac{1}{2}(d-1)+\mathrm{i} \rho\right) \Gamma\left(\frac{1}{2}(d-1)-\mathrm{i} \rho\right) \\
& \times\left[\left(p_{12}\right)^{2}-1\right]^{-1(d-2)} P_{-\frac{1}{2}+\mathrm{i} \rho}^{-\frac{1}{4}(d-2)}\left(-p_{12}(\epsilon)\right) . \tag{9.11}
\end{align*}
$$

Furthermore, formula (11.16) on $p 1713$ should be replaced by

$$
\begin{gather*}
-2 \lambda_{\mathrm{R}}=-2 \lambda+\frac{1}{2}\left(\frac{m^{2}}{4 \pi}\right)^{\omega} \frac{J_{0}(\omega)}{\omega}  \tag{11.16}\\
-2 \omega(2 \omega-1)\left(16 \pi G_{\mathrm{R}}\right)^{-1}=-2 \omega(\omega-1)(16 \pi G)^{-1}+\frac{1}{2 m^{2}}\left(\frac{m^{2}}{4 \pi}\right) \frac{\omega J_{1}(\omega)}{\omega-1}
\end{gather*}
$$

and on p 1717 formula (A.29) by

$$
\begin{equation*}
\tilde{\rho}_{\mathrm{I}}(\dot{g})=\stackrel{g}{g}^{\prime}=\tilde{\rho}_{\mathrm{PT}}^{\prime}(\dot{g}), \quad \tilde{\rho}_{\mathrm{P}}^{\prime}(\dot{g})=\dot{g}^{\dagger-1}=\tilde{\rho}_{\mathrm{T}}^{\prime}(\dot{g}) . \tag{A.29}
\end{equation*}
$$

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Sorking R 1977 J. Phys. A: Math. Gen. 10 717-25
The third sentence of the second footnote on page 722 should be deleted. Since in § 4 $\mathscr{F}^{\mu \nu}$ is being treated as axial, the induced $H$-tensor $\mathbb{E}$ is well defined (as an axial vector density) independently of any external orientation for $H$. (Proof. It is the $H$-dual of the pullback to $H$ of the $M$-dual of $\mathscr{F}^{\mu \nu}$.) It is only for polar $\mathscr{F}^{\mu \nu}$ that $\mathbb{E}$ is what one might call 'externally axial'.

